

EoMs for Translation

$$\frac{d(m\vec{V})}{dt}_{abs} = \sum \vec{F}_{ext}$$

$$\vec{a}_{b_{abs}} = \left. \frac{d\vec{V}_b}{dt} \right|_{Oxyz} + \vec{\omega}_b \times \vec{V}_b$$

$$\vec{V}_b = \begin{bmatrix} U \\ V \\ W \end{bmatrix} \quad \vec{\omega}_b = \begin{bmatrix} P \\ Q \\ R \end{bmatrix}$$

so we have

$$\begin{aligned} \vec{a}_{b_{abs}} &= \begin{bmatrix} \dot{U} \\ \dot{V} \\ \dot{W} \end{bmatrix} + \begin{vmatrix} i & j & k \\ P & Q & R \\ U & V & W \end{vmatrix} \\ &= \begin{bmatrix} \dot{U} + QW - RV \\ \dot{V} + RU - PW \\ \dot{W} + PV - QU \end{bmatrix} \end{aligned}$$

we have defined our **absolute acceleration terms in body axes**, which means we can define the LHS of Equation

$$\frac{d(m\vec{V})}{dt}_{abs} = m \begin{bmatrix} \dot{U} + QW - RV \\ \dot{V} + RU - PW \\ \dot{W} + PV - QU \end{bmatrix}$$

$$\sum \vec{F}_b = \vec{F}_{G_b} + \vec{F}_{A_b} + \vec{F}_{P_b}$$

We presume the aircraft has one or more propulsions providing thrust, T , along a vector defined in the $\mathbf{x/z}$ plane, at angle θ_T to \mathbf{x} . We include an additional term for side wash effects due to propulsion, \mathbf{F}_{Ty}

$$\vec{F}_{T_b} = \begin{bmatrix} F_{Tx} \\ F_{Ty} \\ F_{Tz} \end{bmatrix} = \begin{bmatrix} T \cdot \cos \theta_T \\ F_{Ty} \\ -T \cdot \sin \theta_T \end{bmatrix}$$

equations of motion for translation:

$$m \begin{bmatrix} \dot{U} + QW - RV \\ \dot{V} + RU - PW \\ \dot{W} + PV - QU \end{bmatrix} = \begin{bmatrix} -mg \sin \theta - D \cos \alpha + L \sin \alpha + T \cos \theta_T \\ mg \sin \phi \cos \theta + F_{Ay} + F_{Ty} \\ mg \cos \phi \cos \theta - D \sin \alpha - L \cos \alpha - T \sin \theta_T \end{bmatrix}$$

EoMs for Rotation

Newton's second law: *the angular acceleration is proportional to the net torque and inversely proportional to the moment of inertia. angular momentum is defined as:*

$$\vec{H} = \vec{r} \times (m\vec{V}) \quad \left. \frac{d\vec{H}}{dt} \right|_{abs} = \vec{M}$$

$$\begin{aligned} \vec{H} &= \begin{bmatrix} H_x \\ H_y \\ H_z \end{bmatrix} = \begin{bmatrix} P \cdot I_{xx} - Q \cdot I_{xy} - R \cdot I_{xz} \\ Q \cdot I_{yy} - R \cdot I_{yz} - P \cdot I_{xy} \\ R \cdot I_{zz} - P \cdot I_{xz} - Q \cdot I_{yz} \end{bmatrix} \\ &= \begin{bmatrix} I_{xx} & -I_{xy} & -I_{xz} \\ -I_{yx} & I_{yy} & -I_{yz} \\ -I_{zx} & -I_{zy} & I_{zz} \end{bmatrix} \begin{bmatrix} P \\ Q \\ R \end{bmatrix} \\ &= [I] \vec{\omega} \end{aligned}$$

Simplification of the Inertia Tensor via symmetry

Most aircraft are symmetric about a longitudinal/vertical plane, and will also tend to have the principal yaw tensor axis aligned with Zb. This means we can assume that: $\mathbf{I}_{xy} = \mathbf{I}_{yx} = \mathbf{I}_{zy} = \mathbf{I}_{yz} = 0$

$$\vec{H} = \begin{bmatrix} I_{xx} & 0 & -I_{xz} \\ 0 & I_{yy} & 0 \\ -I_{xz} & 0 & I_{zz} \end{bmatrix} \begin{bmatrix} P \\ Q \\ R \end{bmatrix}$$

$$= [I] \vec{\omega}$$

$$\begin{bmatrix} H_x \\ H_y \\ H_z \end{bmatrix} = \begin{bmatrix} I_{xx} \cdot P - I_{xz} \cdot R \\ I_{yy} \cdot Q \\ -I_{xz} \cdot P + I_{zz} \cdot R \end{bmatrix}$$

Absolute rate of change of angular momentum

$$\begin{aligned} \frac{d\vec{H}}{dt}_{abs} &= \frac{d\vec{H}}{dt}_{Oxyz} + \vec{\omega} \times \vec{H}_{Oxyz} \\ &= \frac{d[I] \vec{\omega}}{dt}_{Oxyz} + \vec{\omega} \times ([I] \vec{\omega}) \\ &= \frac{d[I]}{dt} \vec{\omega} + [I] \frac{d\vec{\omega}}{dt} + \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ P & Q & R \\ (I_{xx} \cdot P - I_{xz} \cdot R) & (I_{yy} \cdot Q) & (-I_{xz} \cdot P + I_{zz} \cdot R) \end{vmatrix} \end{aligned}$$

The products and moments of inertia are constant in time, hence the time rate of change of [I] cancels to zero

$$\begin{aligned} \frac{d\vec{H}}{dt}_{abs} &= \frac{d[I]}{dt} \cdot \vec{\omega} + \frac{d\vec{\omega}}{dt} \cdot [I] + \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ P & Q & R \\ (I_{xx} \cdot P - I_{xz} \cdot R) & (I_{yy} \cdot Q) & (-I_{xz} \cdot P + I_{zz} \cdot R) \end{vmatrix} \\ &= \begin{bmatrix} I_{xx} \cdot \dot{P} - I_{xz} \cdot \dot{R} \\ I_{yy} \cdot \dot{Q} \\ I_{zz} \cdot \dot{R} - I_{xz} \cdot \dot{P} \end{bmatrix} + \begin{bmatrix} Q(I_{zz} \cdot R - I_{xz} \cdot P) - R(I_{yy} \cdot Q) \\ R(I_{xx} \cdot P - I_{xz} \cdot R) - P(I_{zz} \cdot R - I_{xx} \cdot P) \\ P(I_{yy} \cdot Q) - Q(I_{xx} \cdot P - I_{xz} \cdot R) \end{bmatrix} \end{aligned}$$

collecting terms of the moments and products of inertia yields

$$\begin{bmatrix} \dot{H}_x \\ \dot{H}_y \\ \dot{H}_z \end{bmatrix}_{abs} = \begin{bmatrix} \dot{P} \cdot I_{xx} & + & Q \cdot R (I_{zz} - I_{yy}) & - & (\dot{R} + P \cdot Q) I_{xz} \\ \dot{Q} \cdot I_{yy} & + & P \cdot R (I_{xx} - I_{zz}) & + & (P^2 - R^2) I_{xz} \\ \dot{R} \cdot I_{zz} & + & P \cdot Q (I_{yy} - I_{xx}) & + & (Q \cdot R - \dot{P}) I_{xz} \end{bmatrix} = \begin{bmatrix} L \\ M \\ N \end{bmatrix}$$

Summary of the equations of motion#

With the translational (Eq. (29)) and rotational (Eq. (33)) defined, and the relationship between the Euler angles and the body rates (Eqs. eq:eulerratetoomega and (18)), fully unconstrained flight in 6DoF can be described.

In total twelve equations have been derived (body/attitude rates are only three but the matrix is tough to invert, so I'll include both versions below) that can describe aircraft position and attitude in a Newtonian framework:

Translational Motion:

$$m \begin{bmatrix} \dot{U} + QW - RV \\ \dot{V} + RU - PW \\ \dot{W} + PV - QU \end{bmatrix} = \begin{bmatrix} -mg \sin \theta - D \cos \alpha + L \sin \alpha + T \cos \theta_T \\ mg \sin \phi \cos \theta + F_{A_Y} + F_{T_Y} \\ mq \cos \phi \cos \theta - D \sin \alpha - L \cos \alpha - T \sin \theta_T \end{bmatrix}$$

Body angular rate due to an attitude rate:

$$\begin{bmatrix} P \\ Q \\ R \end{bmatrix} = \begin{bmatrix} 1 & 0 & -\sin \theta \\ 0 & \cos \phi & \sin \phi \cos \theta \\ 0 & -\sin \phi & \cos \phi \cos \theta \end{bmatrix} \begin{bmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{bmatrix}$$

Angular Motion:

$$\begin{bmatrix} \dot{P} \cdot I_{xx} + Q \cdot R (I_{zz} - I_{yy}) - (\dot{R} + P \cdot Q) I_{xz} \\ \dot{Q} \cdot I_{yy} + P \cdot R (I_{xx} - I_{zz}) + (P^2 - R^2) I_{xz} \\ \dot{R} \cdot I_{zz} + P \cdot Q (I_{yy} - I_{xx}) + (Q \cdot R - \dot{P}) I_{xz} \end{bmatrix} = \begin{bmatrix} L \\ M \\ N \end{bmatrix}$$

Attitude rate due to a body rate

$$\begin{bmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{bmatrix} = \begin{bmatrix} 1 & \sin \phi \tan \theta & \cos \phi \tan \theta \\ 0 & \cos \phi & -\sin \phi \\ 0 & \frac{\sin \phi}{\cos \theta} & \frac{\cos \phi}{\cos \theta} \end{bmatrix} \begin{bmatrix} P \\ Q \\ R \end{bmatrix}$$

Stability Derivatives

Reduced derivative force and moment perturbations:

$$\Delta X = \left. \frac{\partial X}{\partial u} \right|_0 u + \left. \frac{\partial X}{\partial w} \right|_0 w$$

$$\Delta Y = \left. \frac{\partial Y}{\partial v} \right|_0 v + \left. \frac{\partial Y}{\partial \delta_r} \right|_0 \delta_r$$

$$\Delta Z = \left. \frac{\partial Z}{\partial u} \right|_0 u + \left. \frac{\partial Z}{\partial w} \right|_0 w + \left. \frac{\partial Z}{\partial \delta_e} \right|_0 \delta_e$$

$$\Delta L = \left. \frac{\partial L}{\partial v} \right|_0 v + \left. \frac{\partial L}{\partial p} \right|_0 p + \left. \frac{\partial L}{\partial r} \right|_0 r + \left. \frac{\partial L}{\partial \delta_r} \right|_0 \delta_r + \left. \frac{\partial L}{\partial \delta_a} \right|_0 \delta_a$$

$$\Delta M = \left. \frac{\partial M}{\partial u} \right|_0 u + \left. \frac{\partial M}{\partial w} \right|_0 w + \left. \frac{\partial M}{\partial \dot{w}} \right|_0 \dot{w} + \left. \frac{\partial M}{\partial q} \right|_0 q + \left. \frac{\partial M}{\partial \delta_a} \right|_0 \delta_a$$

$$\Delta N = \left. \frac{\partial N}{\partial v} \right|_0 v + \left. \frac{\partial N}{\partial p} \right|_0 p + \left. \frac{\partial N}{\partial r} \right|_0 r + \left. \frac{\partial N}{\partial \delta_r} \right|_0 \delta_r + \left. \frac{\partial N}{\partial \delta_a} \right|_0 \delta_a$$

The **linear force and moment perturbations**, may be substituted into the **linear equations of motion**, to have:

$$m\dot{u} = \left. \frac{\partial X}{\partial u} \right|_0 u + \left. \frac{\partial X}{\partial w} \right|_0 w - mg \cdot \cos \theta_0 \cdot \theta'$$

$$m\dot{v} = \left. \frac{\partial Y}{\partial v} \right|_0 v - mU_0 r + mg \cdot \cos \theta_0 \cdot \phi' \left. \frac{\partial Y}{\partial \delta_r} \right|_0 \delta_r$$

$$m\dot{w} = \left. \frac{\partial Z}{\partial u} \right|_0 u + \left. \frac{\partial Z}{\partial w} \right|_0 w + mU_0 q - mg \cdot \sin \Theta_0 \cdot \theta' + \left. \frac{\partial Z}{\partial \delta_e} \right|_0 \delta_e$$

$$I_{xx}\dot{p} - I_{xz}\dot{r} = \left. \frac{\partial L}{\partial v} \right|_0 v + \left. \frac{\partial L}{\partial p} \right|_0 p + \left. \frac{\partial L}{\partial r} \right|_0 r + \left. \frac{\partial L}{\partial \delta_r} \right|_0 \delta_r + \left. \frac{\partial L}{\partial \delta_a} \right|_0 \delta_a$$

$$I_{yy}\dot{q} = \left. \frac{\partial M}{\partial u} \right|_0 u + \left. \frac{\partial M}{\partial w} \right|_0 w + \left. \frac{\partial M}{\partial \dot{w}} \right|_0 \dot{w} + \left. \frac{\partial M}{\partial q} \right|_0 q + \left. \frac{\partial M}{\partial \delta_e} \right|_0 \delta_e$$

$$I_{zz}\dot{r} - I_{xz}\dot{p} = \left. \frac{\partial N}{\partial v} \right|_0 v + \left. \frac{\partial N}{\partial p} \right|_0 p + \left. \frac{\partial N}{\partial r} \right|_0 r + \left. \frac{\partial N}{\partial \delta_r} \right|_0 \delta_r + \left. \frac{\partial N}{\partial \delta_a} \right|_0 \delta_a$$

Concise Form of Dimensional EoMs

$$m\dot{u} = \left. \frac{\partial X}{\partial u} \right|_0 u + \left. \frac{\partial X}{\partial w} \right|_0 w - mg \cdot \cos \theta_0 \cdot \theta'$$

$$m\dot{w} = \left. \frac{\partial Z}{\partial u} \right|_0 u + \left. \frac{\partial Z}{\partial w} \right|_0 w + mU_0 q - mg \cdot \sin \Theta_0 \cdot \theta' + \left. \frac{\partial Z}{\partial \delta_e} \right|_0 \delta_e$$

$$I_{yy}\dot{q} = \left. \frac{\partial M}{\partial u} \right|_0 u + \left. \frac{\partial M}{\partial w} \right|_0 w + \left. \frac{\partial M}{\partial \dot{w}} \right|_0 \dot{w} + \left. \frac{\partial M}{\partial q} \right|_0 q + \left. \frac{\partial M}{\partial \delta_e} \right|_0 \delta_e$$

The X-Force equation (*forward speed equation*):

$$\dot{u} = X_u u + X_w w - g \cdot \cos \theta_0 \theta' \quad X_u \triangleq \left. \frac{1}{m} \frac{\partial X}{\partial u} \right|_0, X_w \triangleq \left. \frac{1}{m} \frac{\partial X}{\partial w} \right|_0$$

The Z-Force equation (*heave equation*):

$$\dot{w} = Z_u u + Z_w w + U_0 q - g \cdot \sin \theta_0 \cdot \theta' + Z_{\delta_e} \delta_e$$

$$Z_u \triangleq \left. \frac{1}{m} \frac{\partial Z}{\partial u} \right|_0, Z_w \triangleq \left. \frac{1}{m} \frac{\partial Z}{\partial w} \right|_0, Z_{\delta_e} \triangleq \left. \frac{1}{m} \frac{\partial Z}{\partial \delta_e} \right|_0$$

The M-Moment equation (*pitching moment equation*):

$$\dot{q} = M_u^* u + M_w^* w + M_{\dot{w}}^* \dot{w} + M_q^* q + M_{\delta_e}^* \delta_e$$

$$M_u^* \triangleq M_u + M_{\dot{w}} Z_u \quad M_w^* \triangleq M_w + M_{\dot{w}} Z_w \quad M_q^* \triangleq M_q + M_{\dot{w}} U_0$$

$$M_{\delta_e}^* \triangleq M_{\delta_e} + M_{\dot{w}} Z_{\delta_e} \quad M_{\theta}^* \triangleq -M_{\dot{w}} g \sin \theta_0$$

The linearised Euler pitch equation (*pitch rate kinematic equation*):

$$q \text{ is pitch rate, } = \boldsymbol{\theta}'$$

Now Equations can be written in matrix form to give the **linearised equation of longitudinal motion in concise form**:

$$\begin{bmatrix} \dot{u} \\ \dot{w} \\ \dot{q} \\ \dot{\theta} \end{bmatrix} = \begin{bmatrix} X_u & X_w & 0 & -g \cdot \cos \theta_0 \\ Z_u & Z_w & U_0 & -g \cdot \sin \theta_0 \\ M_u^* & M_w^* & M_q^* & M_{\theta}^* \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} u \\ w \\ q \\ \theta \end{bmatrix} + \begin{bmatrix} 0 \\ Z_{\delta_e} \\ M_{\delta_e}^* \\ 0 \end{bmatrix} [\delta_e]$$

Lateral/Directional EoMs

The asymmetric or lateral/directional EoMs may be similarly manipulated.

$$m\dot{v} = \left. \frac{\partial Y}{\partial v} \right|_0 v - mU_0 r + mg \cdot \cos \theta_0 \cdot \phi' \left. \frac{\partial Y}{\partial \delta_r} \right|_0 \delta_r$$

$$I_{xx}\dot{p} - I_{xz}\dot{r} = \left. \frac{\partial L}{\partial v} \right|_0 v + \left. \frac{\partial L}{\partial p} \right|_0 p + \left. \frac{\partial L}{\partial r} \right|_0 r + \left. \frac{\partial L}{\partial \delta_r} \right|_0 \delta_r + \left. \frac{\partial L}{\partial \delta_a} \right|_0 \delta_a$$

$$I_{zz}\dot{r} - I_{xz}\dot{p} = \left. \frac{\partial N}{\partial v} \right|_0 v + \left. \frac{\partial N}{\partial p} \right|_0 p + \left. \frac{\partial N}{\partial r} \right|_0 r + \left. \frac{\partial N}{\partial \delta_r} \right|_0 \delta_r + \left. \frac{\partial N}{\partial \delta_a} \right|_0 \delta_a$$

The Y-Force equation (*sideslip equation*)

$$\dot{v} = Y_v v - U_e r + g \cdot \cos \theta_0 \phi + Y_{\delta_r} \delta_r \quad \text{where} \quad Y_v \triangleq \frac{1}{m} \frac{\partial Y}{\partial v}, Y_{\delta_r} \triangleq \frac{1}{m} \frac{\partial Y}{\partial \delta_r}$$

The two moment equations are coupled, so need to be decoupled. Dividing Equation by the rolling moment of inertia, **I_{xx}**:

$$\dot{p} = \frac{I_{xz}}{I_{xx}} \dot{r} + L_v v + L_p p + L_r r + L_{\delta_r} \delta_r + L_{\delta_a} \delta_a \quad L_v \triangleq \frac{1}{I_{xx}} \frac{\partial L}{\partial v}, L_p \triangleq \frac{1}{I_{xx}} \frac{\partial L}{\partial p},$$

the same can be performed , dividing by the yawing moment of inertia **I_{zz}**:

$$\dot{r} = \frac{I_{xz}}{I_{zz}} \dot{p} + N_v v + N_p p + N_r r + N_{\delta_r} \delta_r + N_{\delta_a} \delta_a$$

Lastly, the yaw rate and body rate kinematic equations give:

$$r = \dot{\psi} \cos \theta_0 \quad \dot{\psi} = r \cdot \sec \theta_0 \quad p = \dot{\phi} - \dot{\psi} \sin \theta_0 \quad \dot{\phi} = p + r \cdot \tan \theta_0$$

So finally, the *linearised lateral/directional equations of motion* may be expressed in *state-space form*

$$\begin{bmatrix} \dot{v} \\ \dot{p} \\ \dot{r} \\ \dot{\phi} \\ \dot{\psi} \end{bmatrix} = \begin{bmatrix} Y_v & 0 & -U_0 & g \cdot \cos \theta_0 & 0 \\ L_v^* & L_p^* & L_r^* & 0 & 0 \\ N_v^* & N_p^* & N_r^* & 0 & 0 \\ 0 & 1 & \tan \theta_0 & 0 & 0 \\ 0 & 0 & \sec \theta_0 & 0 & 0 \end{bmatrix} \begin{bmatrix} v \\ p \\ r \\ \phi \\ \psi \end{bmatrix} + \begin{bmatrix} Y_{\delta_r} & 0 \\ L_{\delta_r}^* & L_{\delta_a}^* \\ N_{\delta_r}^* & N_{\delta_a}^* \\ 0 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} \delta_r \\ \delta_a \end{bmatrix}$$

The two sets of matrix equations, , are known as *the equations of motion in concise form*.

These are in **state space form**

$$\dot{\vec{x}} = \mathbf{A}\vec{x} + \mathbf{B}\vec{u}$$

where \vec{x} = the state vector (n) \vec{u} = the control matrix (m)

\mathbf{A} = the system matrix (n by n) \mathbf{B} = the control matrix (n by m)